

# Quantum Effects in Josephson Junctions - Part 2

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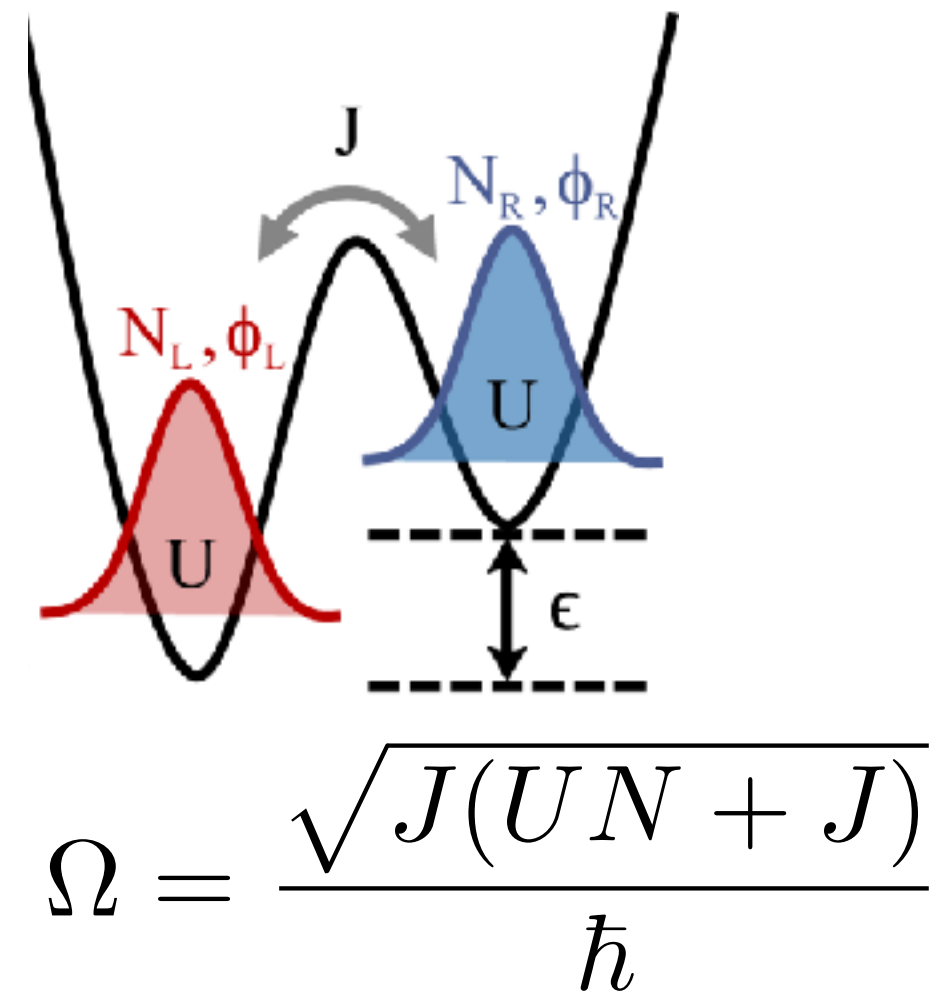
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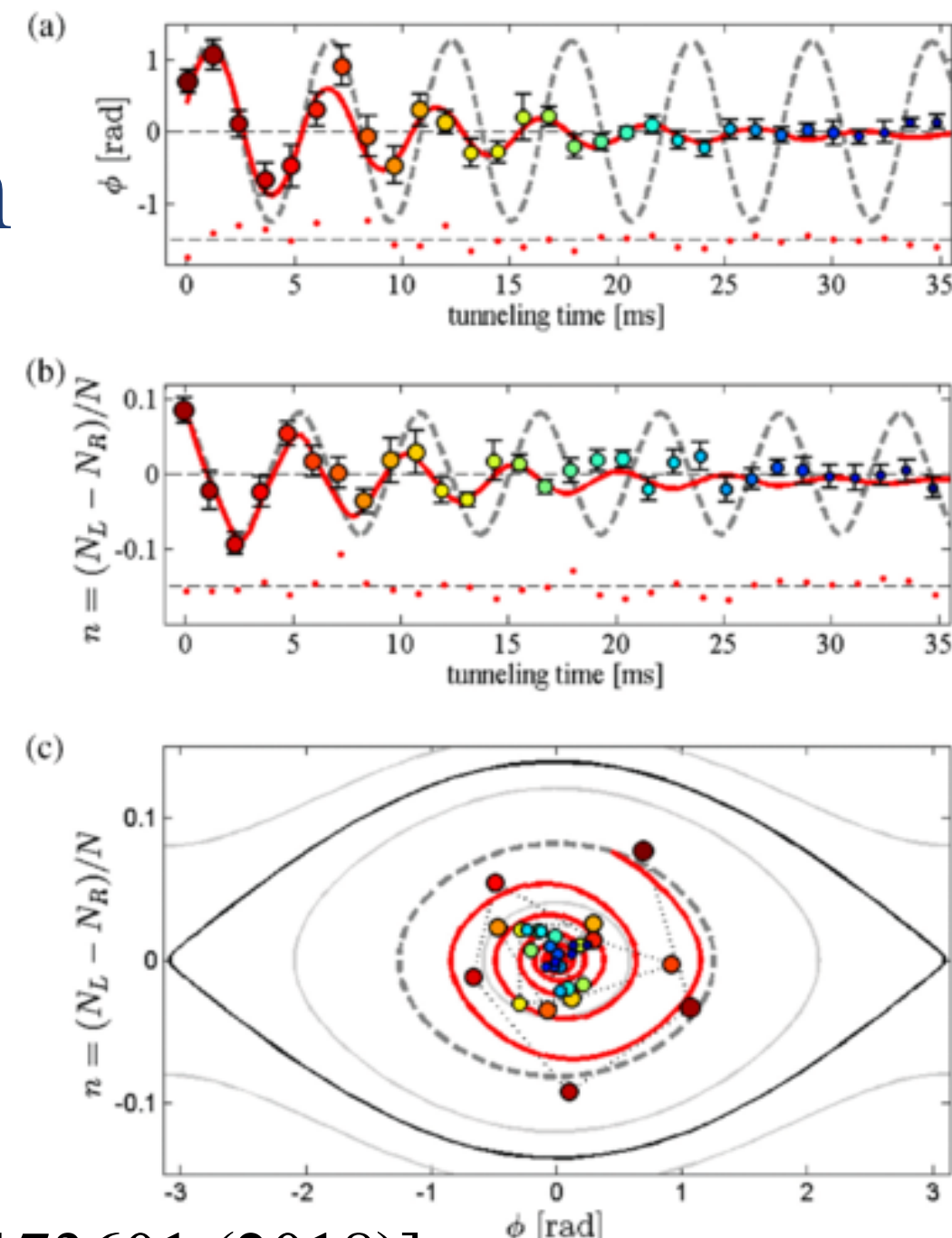


## Introduction

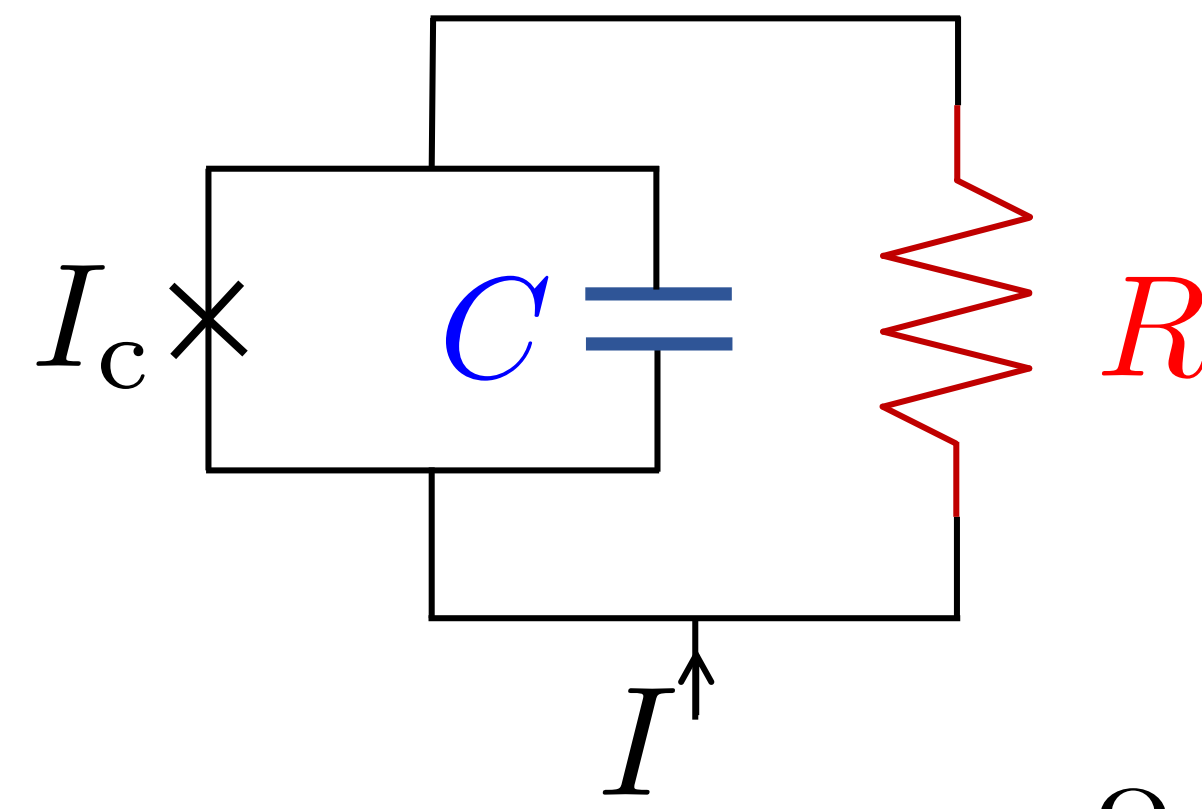
✓ Bose Josephson junction



[M. Pigneur *et al.*, Phys. Rev. Lett. **120**, 173601 (2018)]

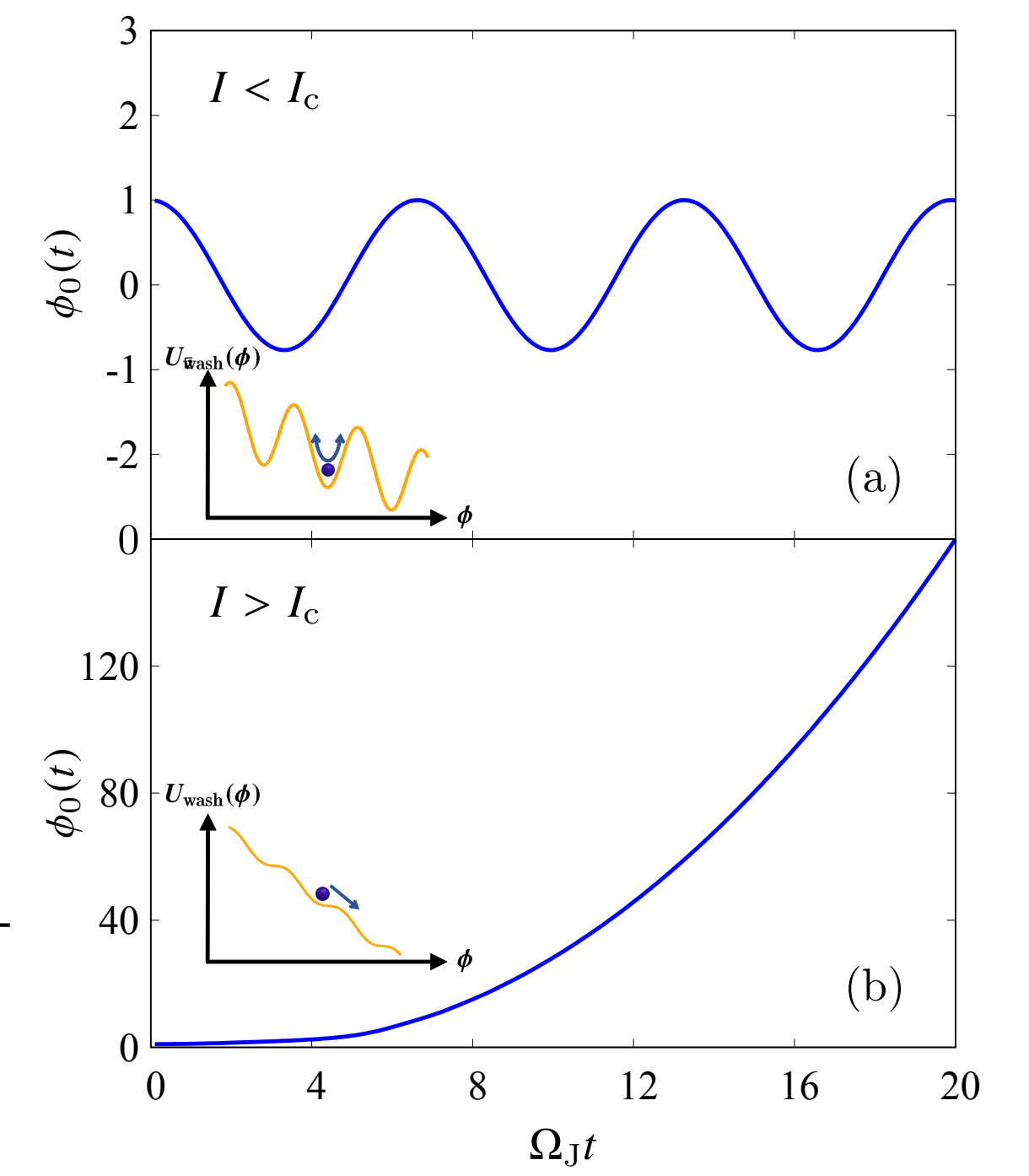


✓ Resistively & capacitively shunted Josephson circuit



$$\Omega = \frac{1}{\hbar} \sqrt{\frac{2e\hbar I_c}{C}}$$

[KF and L. Salasnich, Phys. Rev. B **104**, 014519 (2021)]



Our work

We find the leading quantum correction to Josephson frequency in Josephson junctions.

$$\tilde{\Omega} = \Omega + (\text{Quantum correction})$$

## Bose Josephson junction

$$L = \sum_{j=1,2} \left[ i\hbar\psi_j^* \partial_t \psi_j - \frac{U}{2} |\psi_j|^4 \right] + \frac{J}{2} [\psi_1^* \psi_2 + \psi_2^* \psi_1]$$

$U$  : onsite interaction strength

$J$  : tunnel coupling

$\psi_j(t) = \sqrt{N_j(t)} e^{i\phi_j(t)}$  : complex bosonic field on a site  $j$  at a real-time  $t$

$N_j(t)$  : number of bosons on a site  $j$  at a real-time  $t$

$\phi_j(t)$  : phase angle on a site  $j$  at a real-time  $t$

Total number & total phase  $N(t) = N_1(t) + N_2(t)$   $\bar{\phi}(t) = \phi_1(t) + \phi_2(t)$

Population imbalance & relative phase  $z(t) = \frac{N_1(t) - N_2(t)}{N(t)}$   $\phi(t) = \phi_2(t) - \phi_1(t)$

Integrating out  $\bar{\phi}$  and  $N$

$$L_{\text{rel}}[\phi, z] = \frac{N\hbar}{2} z \dot{\phi} - \frac{UN^2}{4} z^2 + \frac{JN}{2} \sqrt{1 - z^2} \cos \phi$$

► EoM

$$\begin{cases} \ddot{\phi} + \Omega^2 \phi = 0 \\ \ddot{z} + \Omega^2 z = 0 \end{cases}$$

► Josephson frequency

$$\Omega = \frac{\sqrt{J(UN + J)}}{\hbar} = \begin{cases} \Omega_R = \frac{J}{\hbar} & (UN/J \ll 1) \\ \Omega_J = \frac{\sqrt{UNJ}}{\hbar} & (UN/J \gg 1) \end{cases}$$

► Path-integral propagator

$$K(\phi_f, T | \phi_i, 0) = \int_{(\phi_i, 0)}^{(\phi_f, T)} \mathcal{D}\phi \int \mathcal{D}z e^{iS_{\text{rel}}[\phi, z]/\hbar} = \int_{(\phi_i, 0)}^{(\phi_f, T)} \mathcal{D}\phi e^{iS_0[\phi]/\hbar}$$

► Phase action

$$S_0[\phi] = \int dt \left[ \frac{N\hbar^2 \dot{\phi}^2}{4(UN + J \cos \phi)} + \frac{JN}{2} \cos \phi \right]$$

$$\simeq S_J[\phi] = \int dt \left[ \frac{M_J}{2} \dot{\phi}^2 - V(\phi) \right]$$

in the Josephson regime  $UN/J \gg 1$

$$M_J \equiv \frac{\hbar^2}{2U} \quad V(\phi) \equiv \frac{JN}{2} (1 - \cos \phi)$$

$$\lambda \equiv -\frac{JN}{48} \quad = \frac{M_J \Omega_J^2}{2} \phi^2 + U(\phi) = \frac{M_J \Omega_J^2}{2} \phi^2 + \lambda \phi^4 + \mathcal{O}(\phi^6)$$

Quantum fluctuations

## Effective action

$$\Gamma_J[\phi] = \int dt \left[ \frac{Z(\phi)}{2} \dot{\phi}^2 - V_{\text{eff}}(\phi) \right]$$

up to the 2nd order derivative expansion

► Validity conditions

$$|\phi| \ll 1, \quad 1 \ll \frac{UN}{J} \ll N^2$$

►  $\hbar$  expansion

$$Z(\phi) = M_J + Z_1(\phi) + Z_2(\phi) + \dots = M_J + Z_1(\phi) + \mathcal{O}(\hbar^2)$$

$$V_{\text{eff}}(\phi) = \frac{M_J \Omega_J^2}{2} \phi^2 + V_{\text{eff}}^{(1)}(\phi) + V_{\text{eff}}^{(2)}(\phi) + \dots$$

$$= \frac{M_J \Omega_J^2}{2} \phi^2 + V_{\text{eff}}^{(1)}(\phi) + \mathcal{O}(\hbar^2)$$

$$Z_1(\phi) = \frac{\hbar}{32M_J^2} \frac{[\partial_\phi^3 U(\phi)]^2}{[\Omega_J^2 + \partial_\phi^2 U(\phi)/M_J]^{5/2}} = \mathcal{O}(\phi^2)$$

$$V_{\text{eff}}^{(1)}(\phi) = \frac{\hbar}{2} \left[ \sqrt{\Omega_J^2 + \frac{1}{M_J} \partial_\phi^2 U(\phi)} - \Omega_J \right] = \frac{3\hbar\lambda}{M_J \Omega_J} \phi^2 + \mathcal{O}(\phi^4)$$

Josephson frequency with 1-loop correction

$$\tilde{\Omega}_J = \Omega_J \sqrt{1 - \frac{1}{2} \sqrt{\frac{U}{JN}}}$$

## Alternative approach

$$\text{EoM: } \ddot{\phi} + \Omega_J^2 \sin \phi = 0$$

$$\phi(t) = \phi_0(t) + \tilde{\phi}(t) \quad \langle \tilde{\phi} \rangle = 0 \quad \langle \tilde{\phi}^2 \rangle = \sqrt{\frac{U}{JN}}$$

Mean-field configuration fluctuations

$$\ddot{\phi}_0 + \Omega_J^2 \left( 1 - \frac{1}{2} \langle \tilde{\phi}^2 \rangle \right) \sin \phi_0 = 0$$

$$= \tilde{\Omega}_J^2$$

... consistent with the result of effective action formalism

## Summary

[KF, J. Tempere, and L. Salasnich, Phys. Rev. B **105**, 134510 (2022)]

- We have obtained a leading quantum correction to the Josephson frequency in a 1D Josephson junction via effective action formalism.
- Our result has been verified also by another approach through the equation of motion.

✓ Bose Josephson junction

[M. Pigneur *et al.*, Phys. Rev. Lett. **120**, 173601 (2018)]

$$\tilde{\Omega}_J = \Omega_J \sqrt{1 - \frac{1}{2} \sqrt{\frac{U}{JN}}} \quad \frac{\Omega_J - \tilde{\Omega}_J}{\Omega_J} \simeq 0.1\% \quad \text{with } N = 2500 \quad \frac{UN}{2J} \sim 10^2$$

✓ RCSJ junction

[M. H. Devoret, J. M. Martinis, and J. Clarke, Phys. Rev. Lett. **55**, 1908 (1985)]

$$\tilde{\Omega}_J = \frac{1}{\hbar} \sqrt{\frac{2e\hbar I_0}{C}} \sqrt{1 - \frac{1}{2} \sqrt{\frac{2e^3}{\hbar C I_0}}} \quad \frac{\Omega_J - \tilde{\Omega}_J}{\Omega_J} \simeq 0.03\% \quad \text{with } \sqrt{\frac{8e^3}{\hbar C I_0}} \simeq 2.3 \times 10^{-3}$$